

# Summer School

# Shanghai Observatory

趙丰 Benjamin Fong Chao  
Institute of Earth Sciences  
Academia Sinica, Taipei, Taiwan

# Outline

- Linear expansion
  - Vector space
  - Fourier analysis
  - Spherical harmonics
  - EOF/PCA
- Normal modes
  - of musical instruments
  - of Earth
- Inverse problems
- Earth's rotation
  - “Astronomical”
  - “Geophysical”
- Gravity and Geomagnetism

# Vector space

- Dimensionality
- Addition
- Null vector
- Scaling / multiplication
- Unit vector
- Inner product
- norm
- Basis
- Projection / component

# Tensor of degree $n$

- Scalar ( $n = 0$ )
- Vector ( $n = 1$ )
- Tensor of  $n = 2$ ; (matrix)
  - stress
  - strain
- Tensor of  $n = 4$ :
  - elasticity / compliance

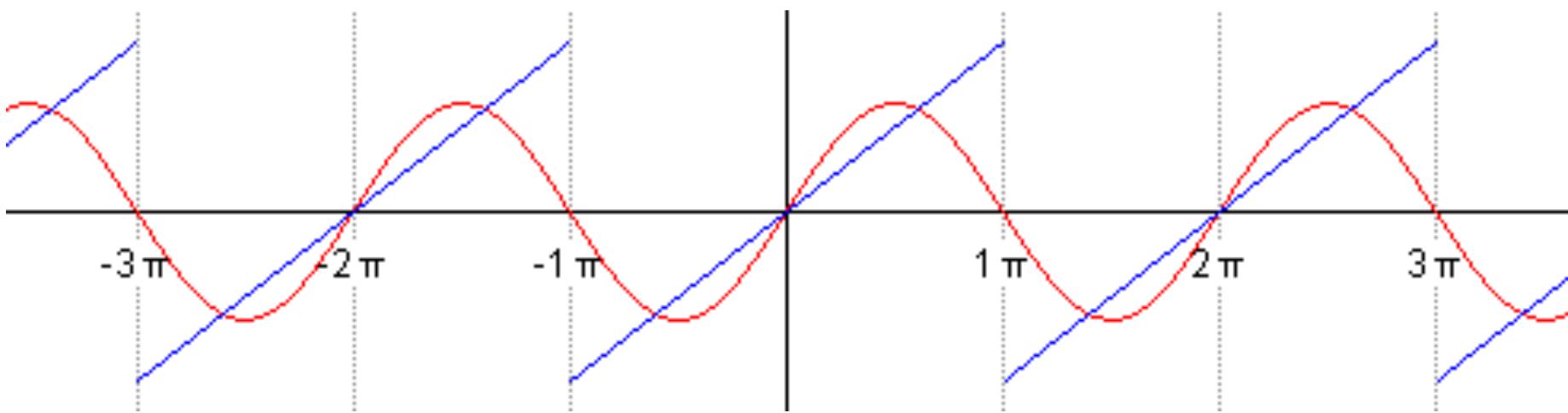
$$\sigma_{ij} = c_{ijkl} \epsilon_{kl}$$

# Hilbert space

- “Function space”
- Infinite dimension
- domain
- Inner product
- Orthogonality
- Basis function
- Completeness

# Fourier analysis

- Basis function = sinusoids
- Cartesian coordinates of dimension  $n$
- Orthogonal
- Complete



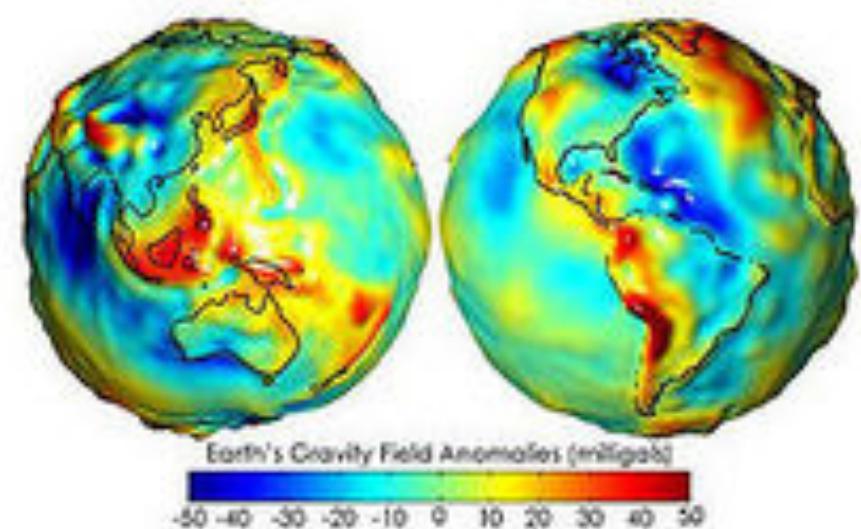
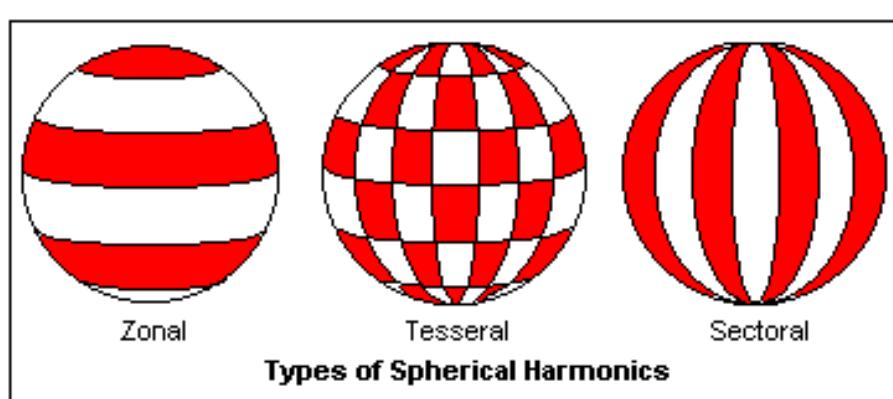
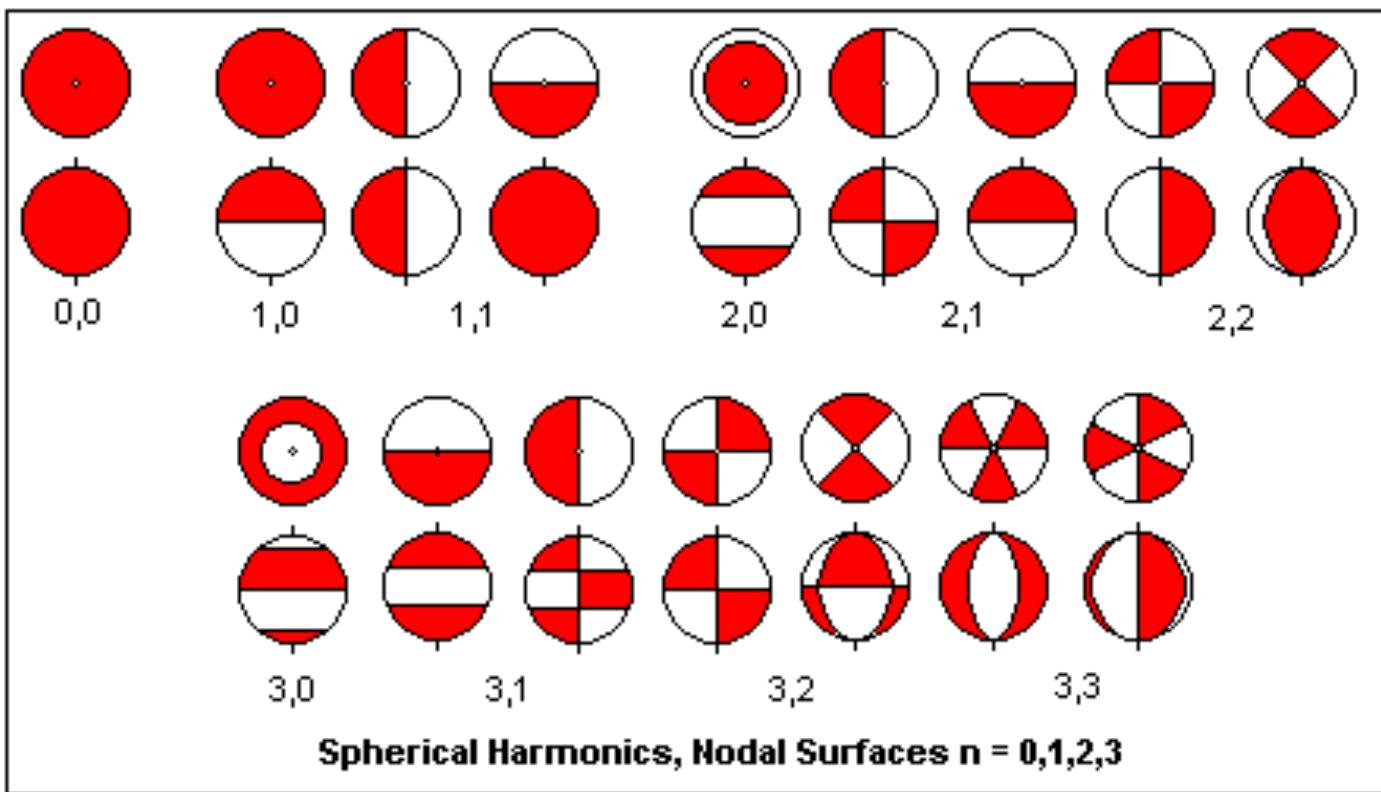
# Spherical Harmonics

- Spherical coordinates
- Satisfying Laplace equation
- Solid harmonics (3-D)
- Surface harmonics (2-D)
- Legendre functions
- Orthogonal
- Complete

$$Y_l^m(\theta, \phi) = \sqrt{\frac{2l+1}{4\pi} \frac{(l-m)!}{(l+m)!}} P_l^m(\cos \theta) e^{im\phi}$$

$$\left\{ \begin{array}{l} Y_{00}=\frac{1}{(4\pi)^{\frac{1}{2}}} \\ Y_{10}=\left(\frac{3}{4\pi}\right)^{\frac{1}{2}}\cos{(\theta)} \\ Y_{1\pm 1}=\mp\left(\frac{3}{8\pi}\right)^{\frac{1}{2}}\sin(\theta)\,e^{\pm i\phi} \\ Y_{20}=\left(\frac{5}{16\pi}\right)^{\frac{1}{2}}(3\cos^2(\theta)-1) \\ Y_{2\pm 1}=\mp\left(\frac{15}{8\pi}\right)^{\frac{1}{2}}\cos{(\theta)}\sin(\theta)\,e^{\pm i\phi} \\ Y_{2\pm 2}=\left(\frac{15}{32\pi}\right)^{\frac{1}{2}}\sin^2(\theta)\,e^{\pm 2i\phi} \end{array} \right.$$

$$\int_0^{2\pi} \int_0^\pi (Y_l^m)^* Y_{l'}^{m'} \sin \theta d\theta d\phi = \delta_{ll'} \delta_{mm'}$$



## EOF (Empirical Orthogonal Function)

- decomposing data matrix into “mode of standing-oscillations”
- presented by spatial pattern and temporal series.

Data matrix  $D(x, t)$ :

$$\begin{bmatrix} x_{11} & x_{12} & \dots & \dots & x_{1p} \\ x_{21} & + & & & \\ \vdots & & & & \vdots \\ \vdots & & & & \vdots \\ x_{n1} & & \dots & \dots & x_{np} \end{bmatrix}$$

observations from all the used stations at time  $j$ .  
 $j = 1, 2, \dots, n$

time series for station  $x_i, i = 1, 2, \dots, p$

Decomposed:  $D(x, t) = \sum_i S_i(x) T_i(t)$ ,

**EOF = the eigen-solutions of the covariance matrix of  $D$  :  $R = D^T D$**

**Spatial pattern  $S_i(x)$  = eigenvectors (orthogonal)**

**Time series  $T_i(t)$  = projection of  $F$  onto the  $i$ -th eigenvector (orthogonal)**

**% variance = eigen-value**

# Normal modes

- Musical instrument
  - Wave equation
  - Boundary condition
  - Propagating wave / normal mode duality
  - 1-D: string
  - 2-D: drum
  - 3-D:
- Earth

# What is music?

Satisfying the wave equation ( $n$ -D) under boundary conditions:

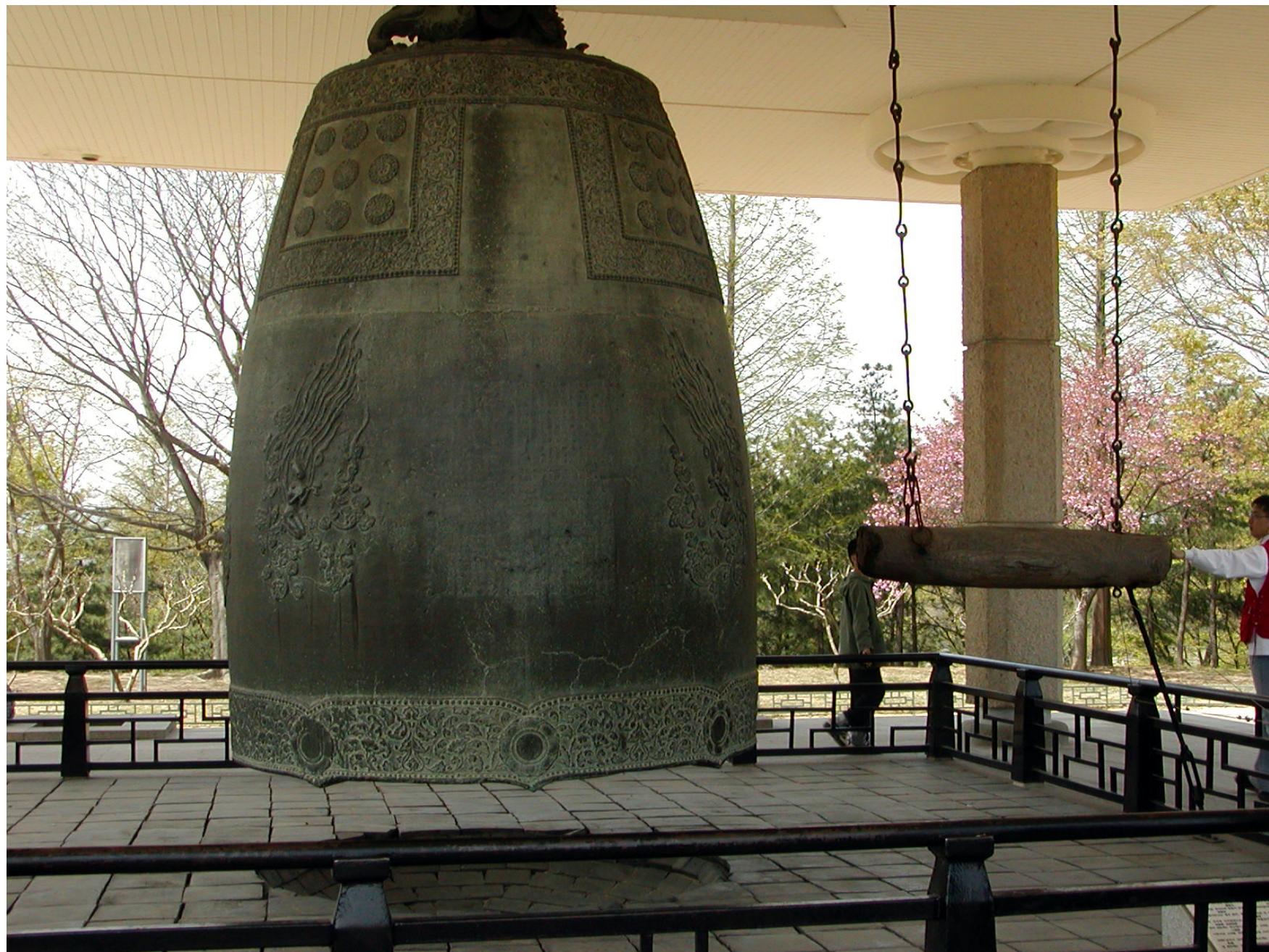
Oscillation of the bell ( $\mathbf{r}, t$ )

=  $\Sigma$  (of all normal modes,  $\mathbb{W}$ )

amplitude (depending on where, how hard you strike, etc., called “excitation”.)

\* normal-mode eigenfunction ( $\mathbf{r}$ ) (depending on the physical property of the bell, e.g, if symmetric, sinusoids in 1-D, Legendre or Bessel functions in 2-D, etc. Earth is 3-D = 2-D + 1-D.)

\* $\exp(i\omega t)$  ( $\omega$  is the normal-mode eigenfrequency, or “natural” resonance frequency = music tones, with imaginary part = natural decay. Quantized because of boundary conditions.)



# A typical seismogram

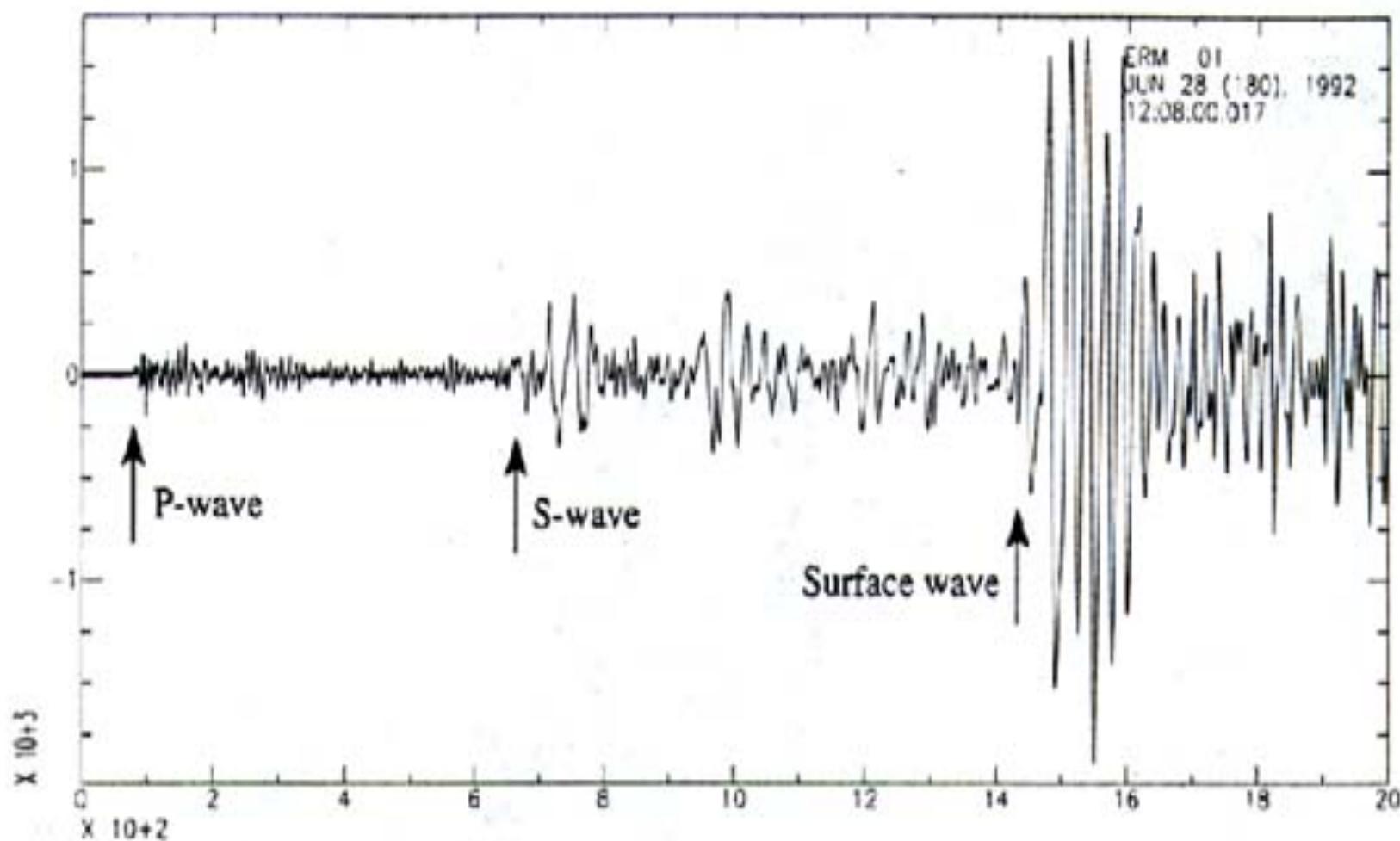
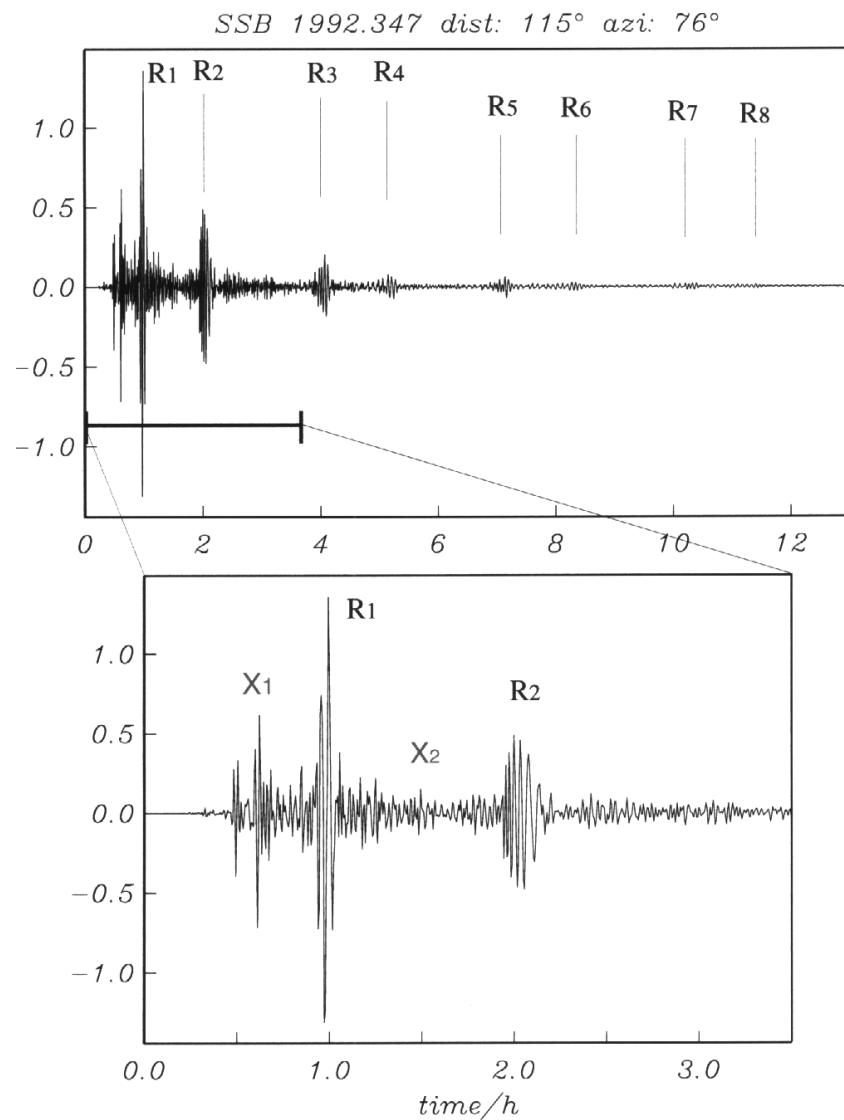
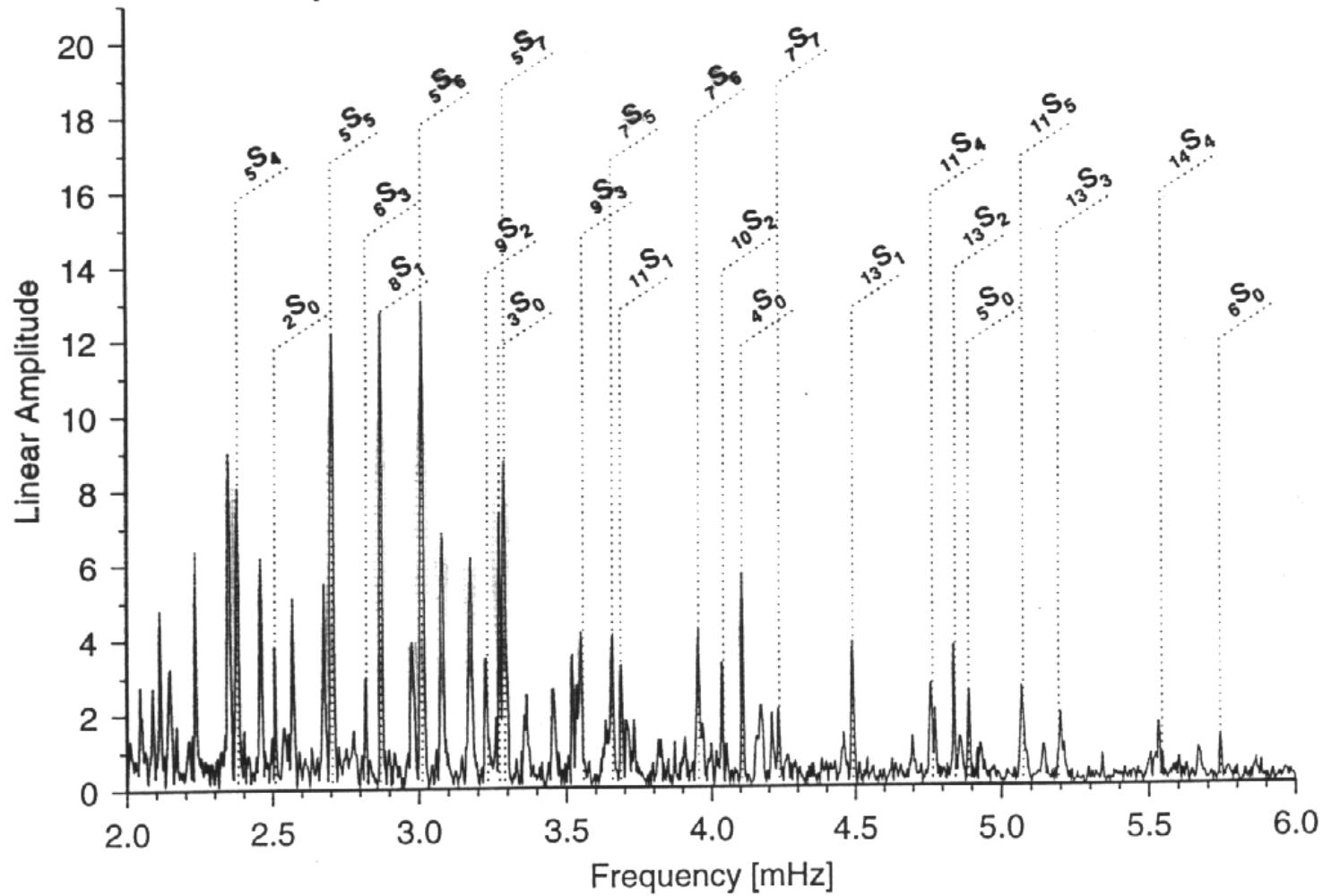


FIGURE 6.4 A seismogram.

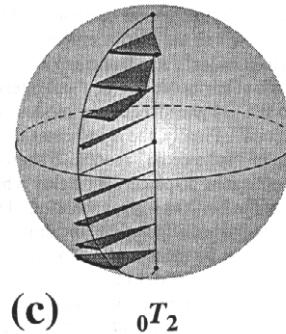
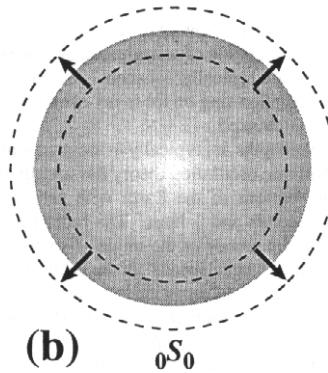
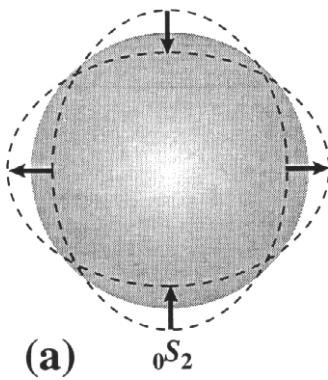


Travelling waves versus standing waves

July 31, 1970 Colombian event recorded at Payson, Arizona

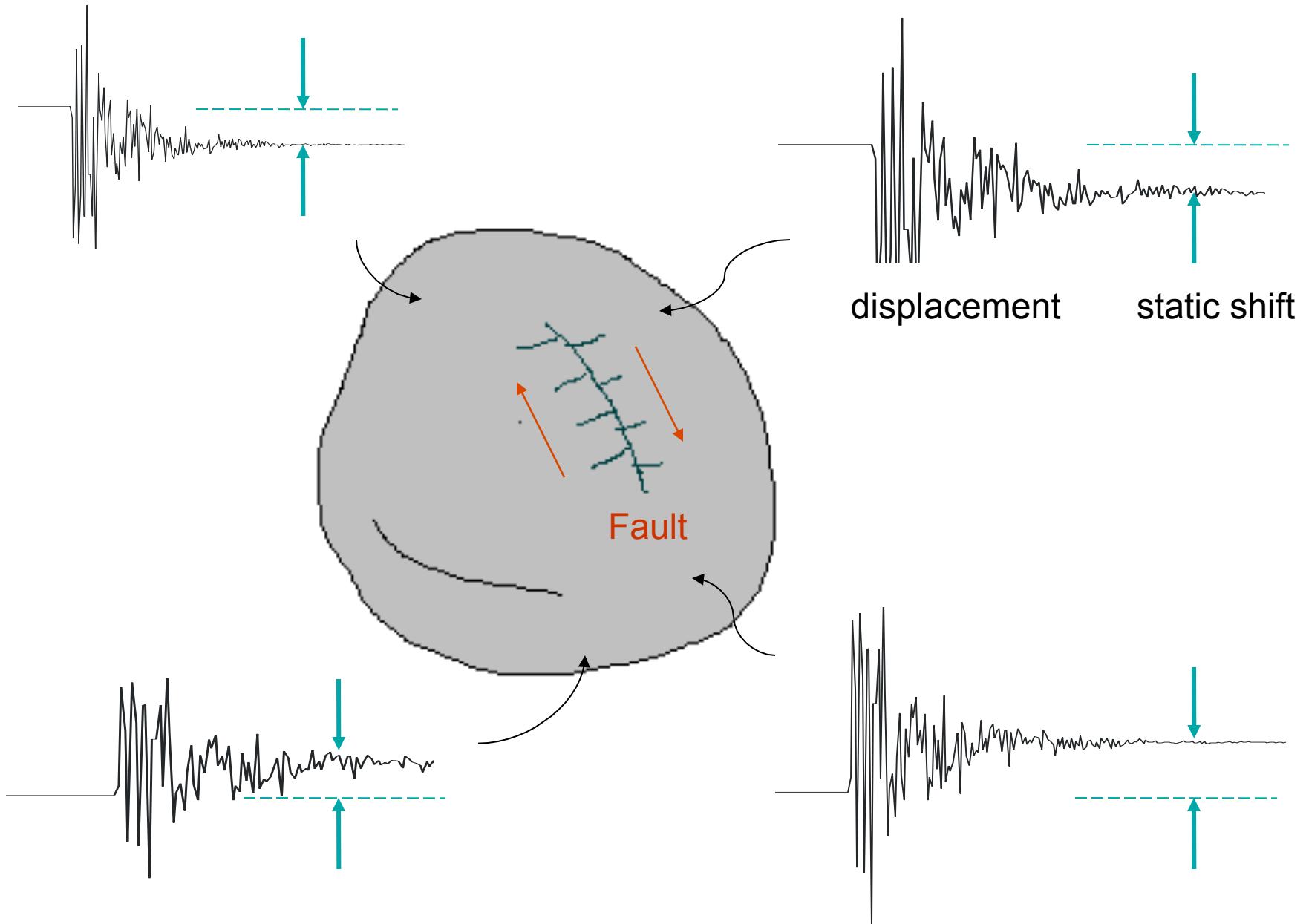


# Different classes of free oscillations



Spheroidal, Radial, and Toroidal

## Static displacement produced by a “fault” in an elastic body



# Earthquake Displacement Field

- Equation of motion

$$\nabla \tau + \mathbf{f}_g + \mathbf{f}_s = \rho(\mathbf{r}) \frac{\partial^2 \mathbf{u}}{\partial t^2}$$

- Solve by expanding displacement field

$$\mathbf{u}(\mathbf{r},t) = \sum_k a_k(t) \mathbf{u}_k^*(\mathbf{r})$$

Normal mode eigenfunctions

$$\mathbf{u}_k(\mathbf{r}) = n U_l(r) Y_l^m(\theta, \phi) \hat{\mathbf{r}} + n V_l(r) \frac{\partial Y_l^m}{\partial \theta} \hat{\boldsymbol{\theta}} + n W_l(r) \frac{1}{\sin \theta} \frac{\partial Y_l^m}{\partial \lambda} \hat{\boldsymbol{\lambda}}$$

Expansion coefficients (note the static limit)

$$a_k(t) = \frac{M_O}{\omega_k^2} \hat{\mathbf{M}} : \mathbf{E}_k^*(\mathbf{r}_s) [\exp(i\omega_k t) + 1]$$

## Co-Seismic Displacement Field

$\mathbf{u}(\mathbf{r}) = \text{oscillations} + \text{static displacement}$

$$= 0 \text{ (as } t \rightarrow \infty) + \sum_{k=0}^{\infty} \omega_k^{-2} \mathbf{u}_k(\mathbf{r}) \mathbf{M} : \mathbf{E}_k^*(\mathbf{r}_f)$$

(Gilbert, 1970 )

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Eigen-mode ( $\mathbf{u}_k(\mathbf{r})$  ,  $\mathbf{E}_k(\mathbf{r}_f)$  ,  $\omega_k$  ,  $k$  = spheroidal and toroidal:  
from SNREI model (e.g., 1066A, B; PREM)

Moment tensor  $\mathbf{M}$ : from Global CMT catalog

# The Scripps gang

## Inverse theory / Normal mode



George Backus



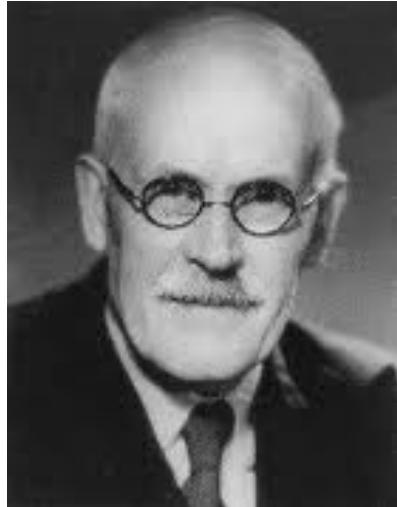
Freeman Gilbert



Robert Parker



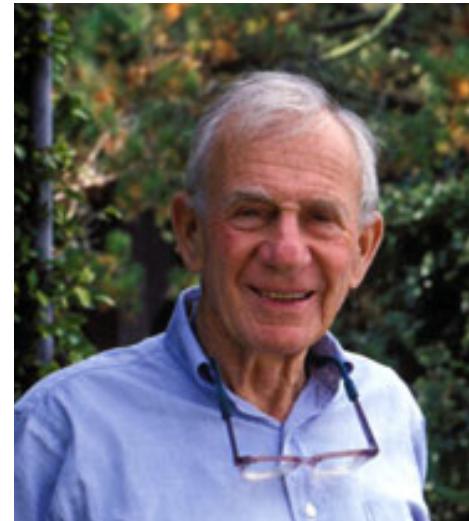
Guy Masters



Sir Harald Jeffreys  
(1891-1989)



Gordon MacDonald  
(1929-2002)



Walter Munk



Kurt Lambeck



Tony Dahlen  
(1942-2007)



John Wahr

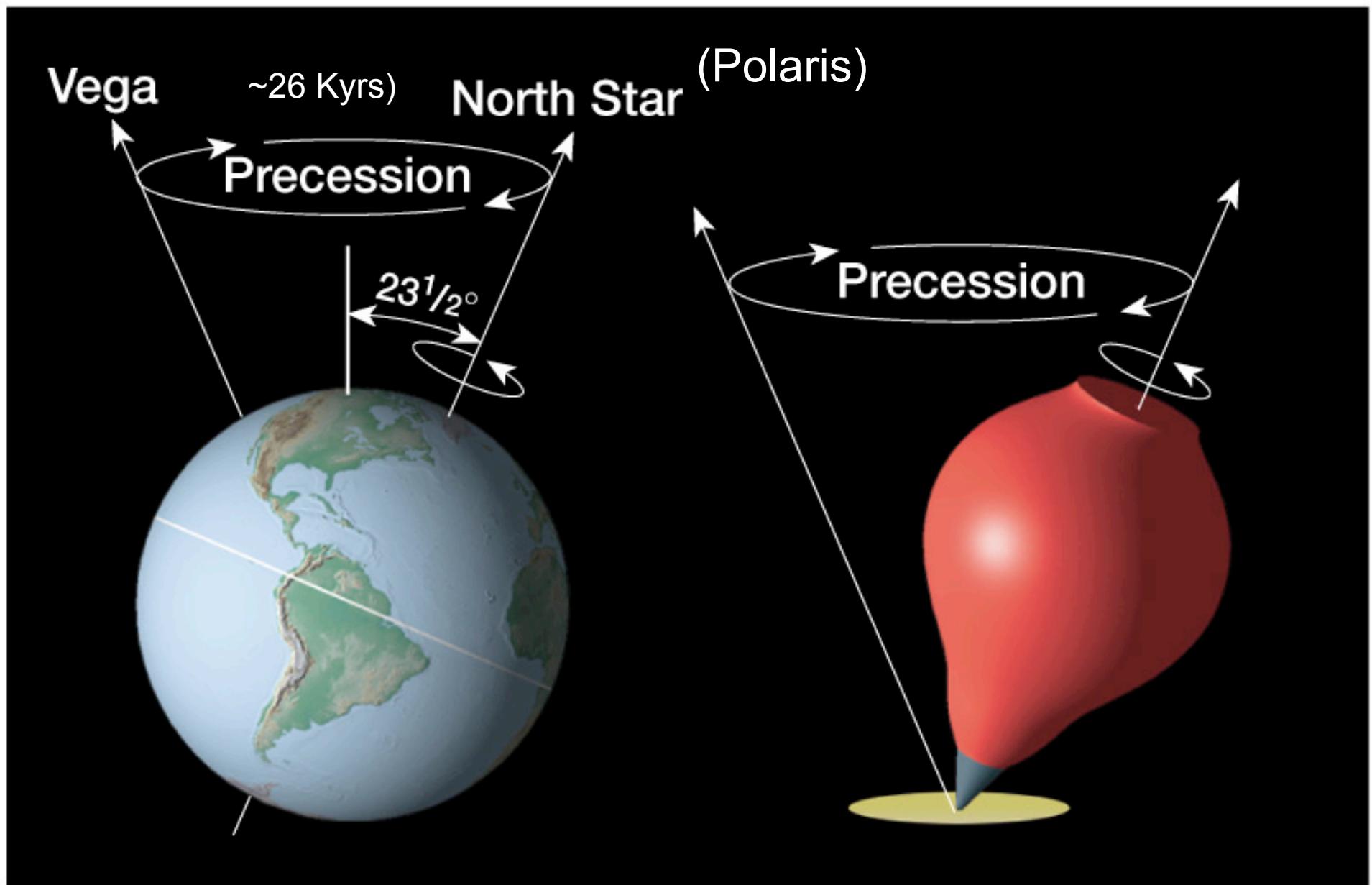
## Earth Rotation

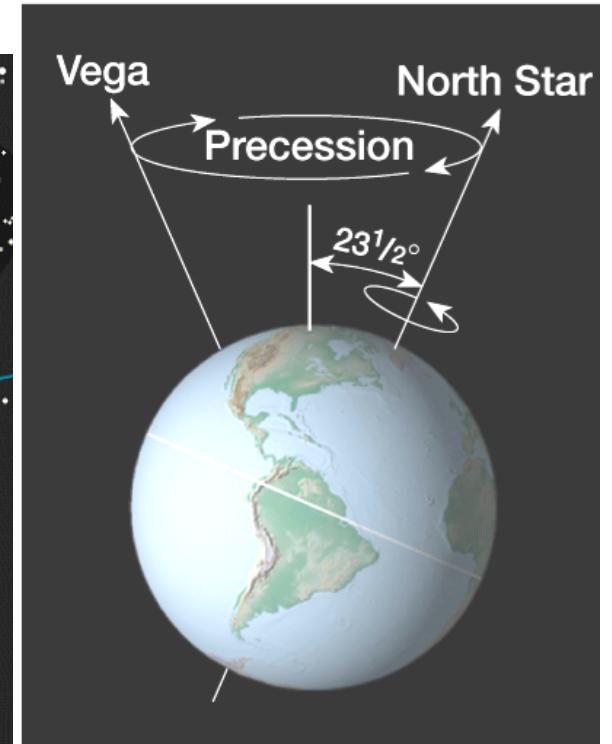
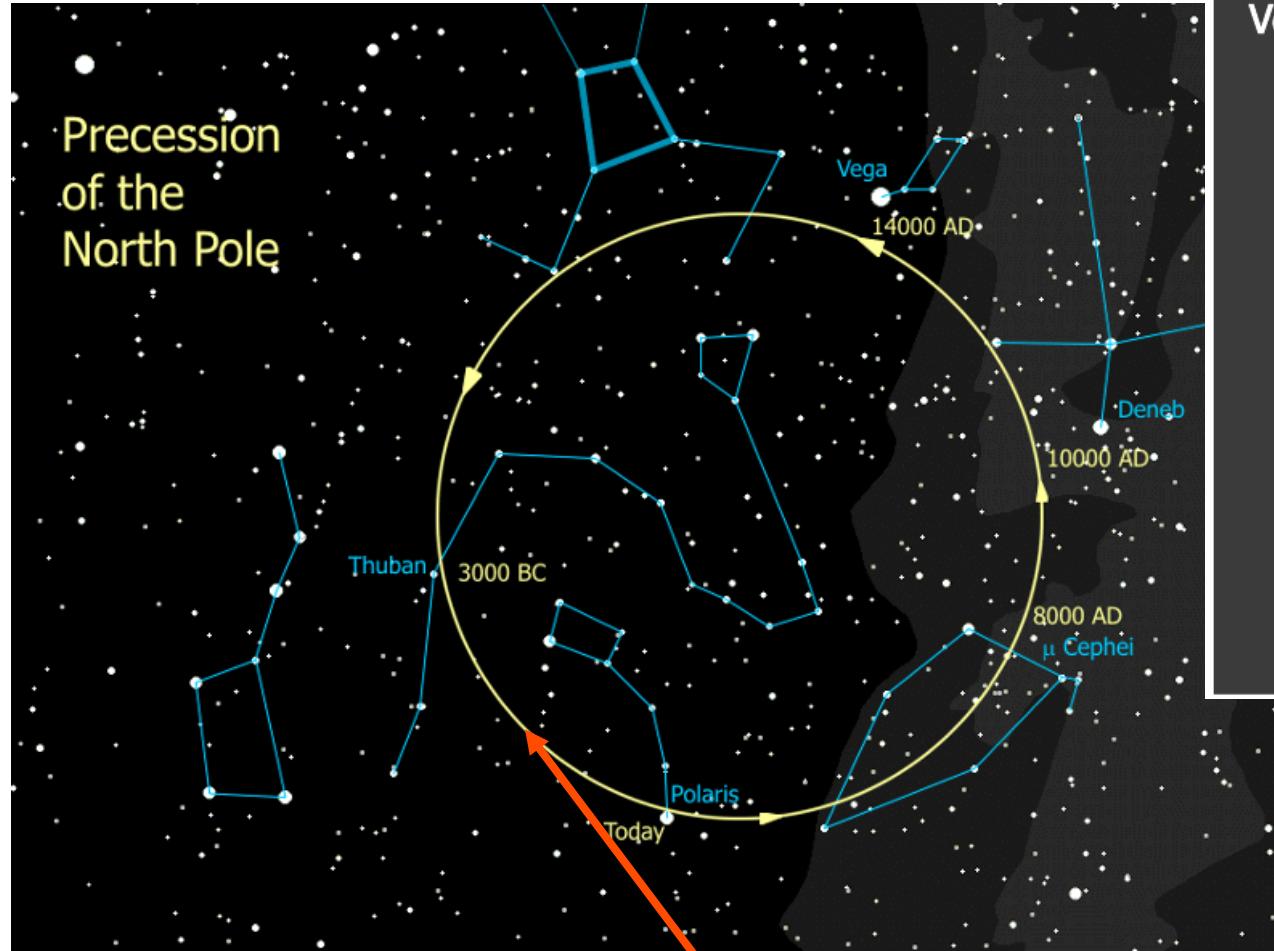
“The Earth precesses/nutates  
like a top.”

“The Earth librates  
like a physical pendulum.”

“The Earth wobbles  
like a frisbee.”

“The Earth precesses/nutates like a top.”





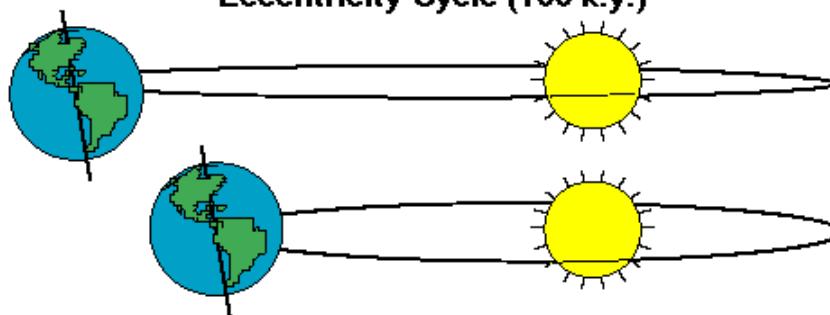
子曰：為政以德，譬如北辰，居其所，而眾星拱之。  
《論語·為政》

**“The Earth wobbles like a frisbee.”**

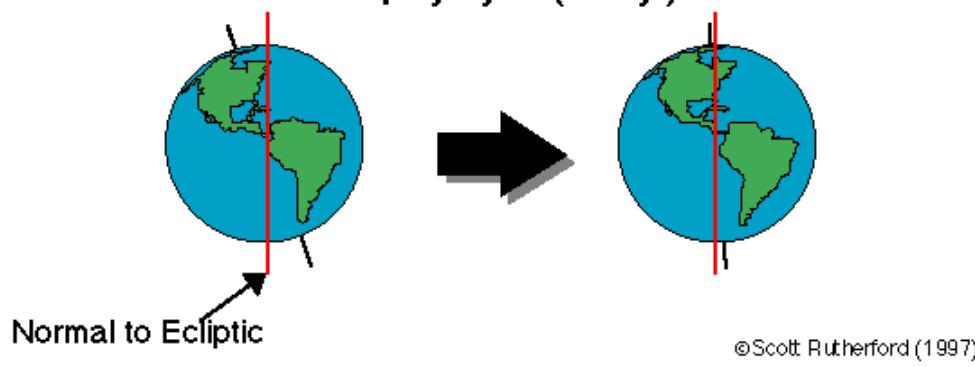


# Milankovitch Cycles

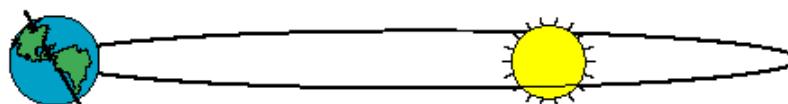
Eccentricity Cycle (100 k.y.)



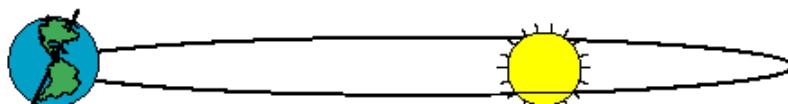
Obliquity Cycle (41 k.y.)



Precession of the Equinoxes (19 and 23 k.y.)

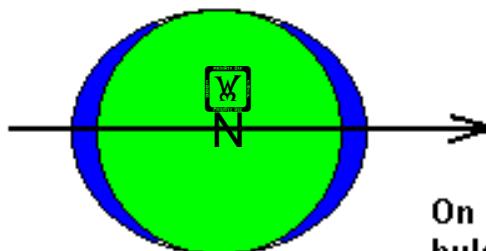


Northern Hemisphere tilted away from the sun at aphelion.

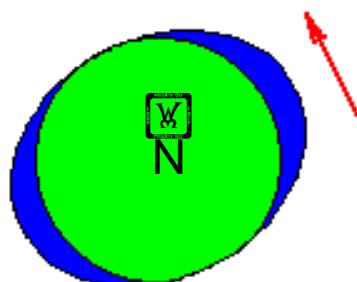


Northern hemisphere tilted toward the sun at aphelion.

## Tidal Braking: Slowing down Earth's rotation and pushing away the Moon

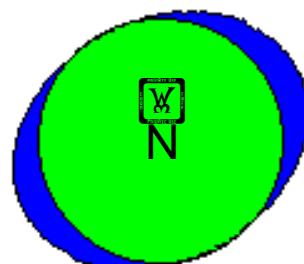


On an ideal, frictionless Earth, tidal bulges would point directly at the Moon.



Moon

On the real Earth, friction and obstructions to ocean movement cause the bulges to lead the Moon.



Moon

The Moon pulls backward on the nearer tidal bulge, slowing Earth's rotation. The nearer bulge, in turn, pulls ahead on the Moon, accelerating it into a higher orbit.

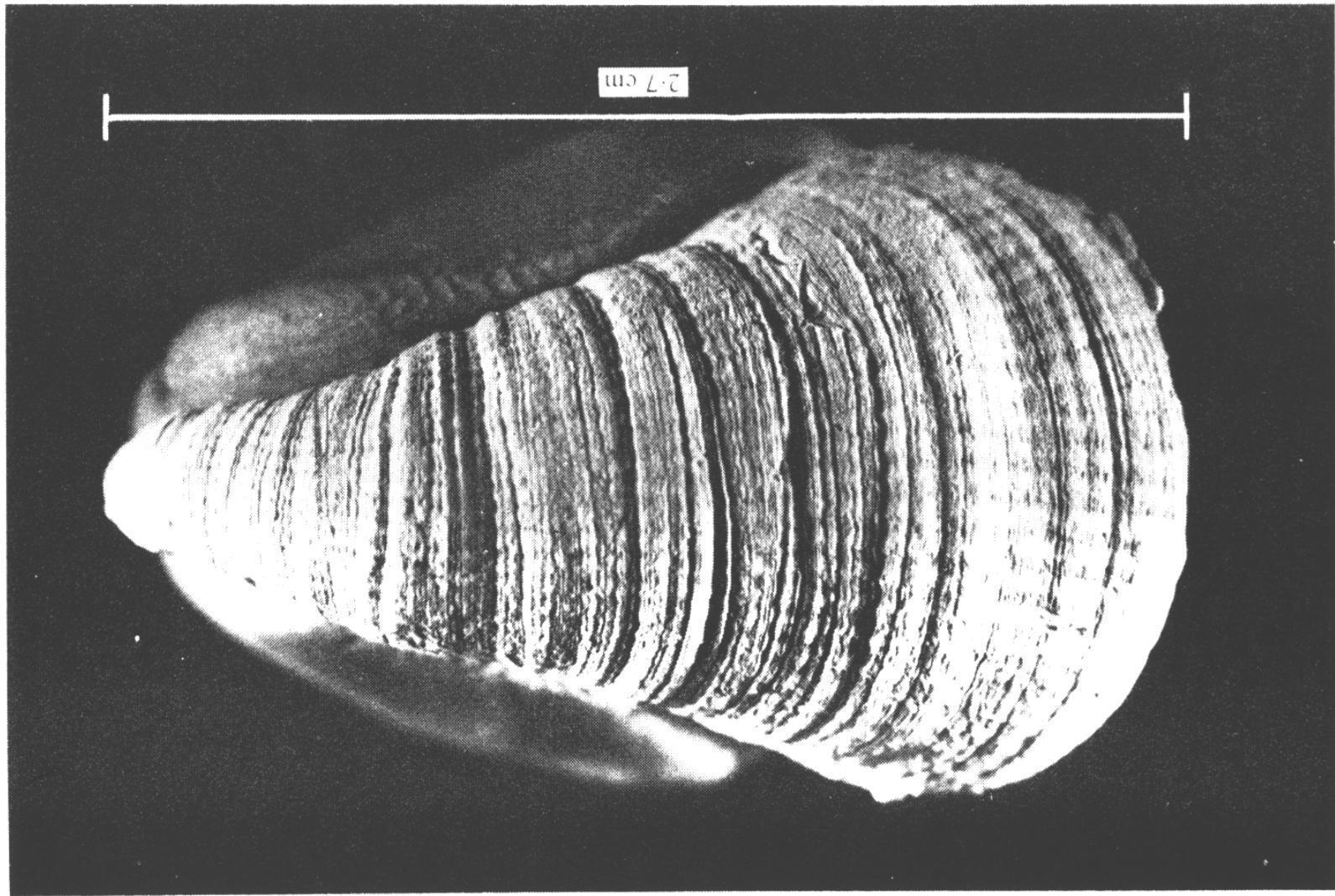
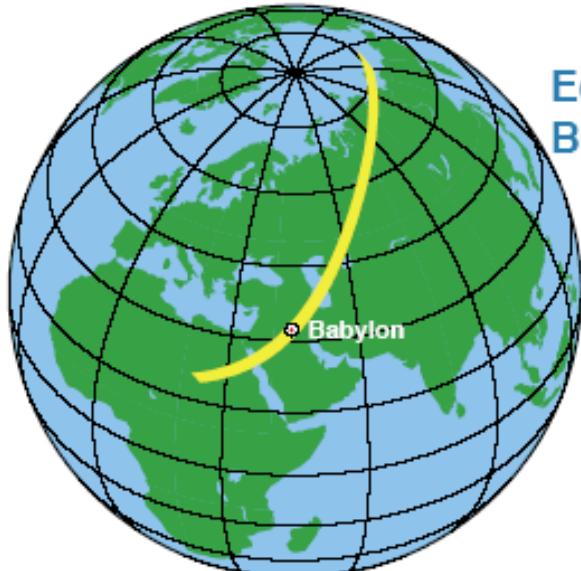


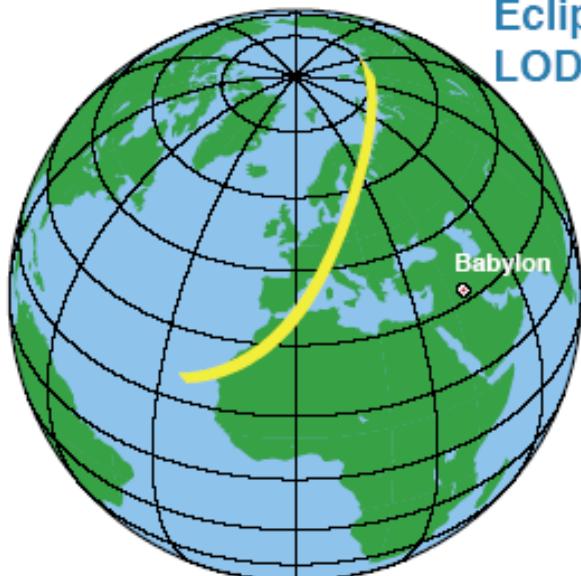
Figure 11.6. Middle Devonian coral epitheca from Michigan, U.S.A., illustrating 13 well-developed bands, each with an average of 30.8 ridges  
*(reproduced by C. T. Sorenson)*

## Secular Braking of Earth's Rotation

### Determination of $d\Omega/dt$ from Ancient Eclipses

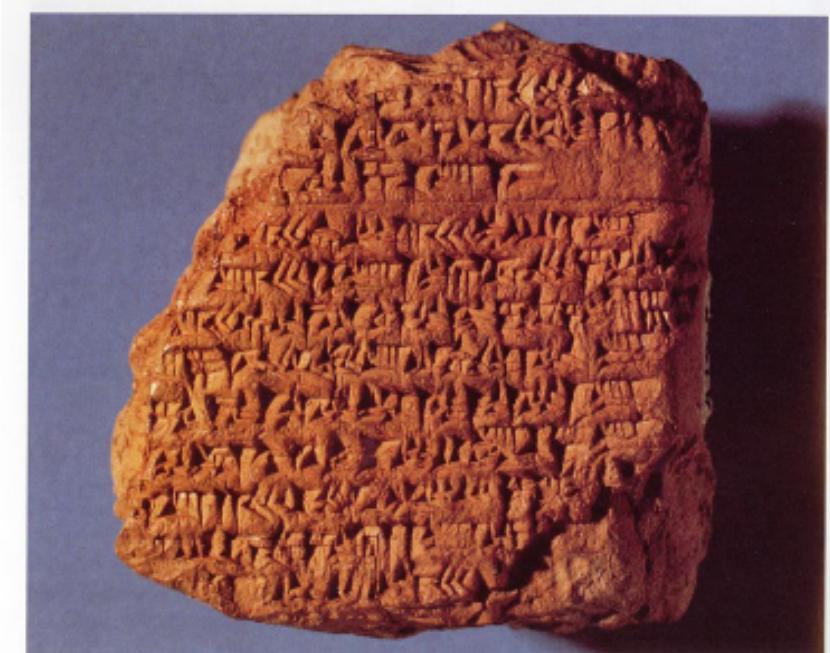


Eclipse path on  
BC 136 April 15



Eclipse path if  
LOD = constant

Eclipse path from Fred Espenak, GSFC



Babylonian diary from the year 67 BC (©The British Museum).

$$\Delta T = 11680 \pm 460 \text{ seconds (3.2 hrs)}$$

(Uncertainties are strict upper/lower bounds)  
(Assumes modern  $dn/dt = -26''/\text{cy}^2$ )

$$\text{Implies } d\Delta/dt = 1.71 \pm 0.07 \text{ ms/century}$$

A Babylonian day was  $\sim 37$  ms shorter than ours.

A very precise estimate from one single eclipse!

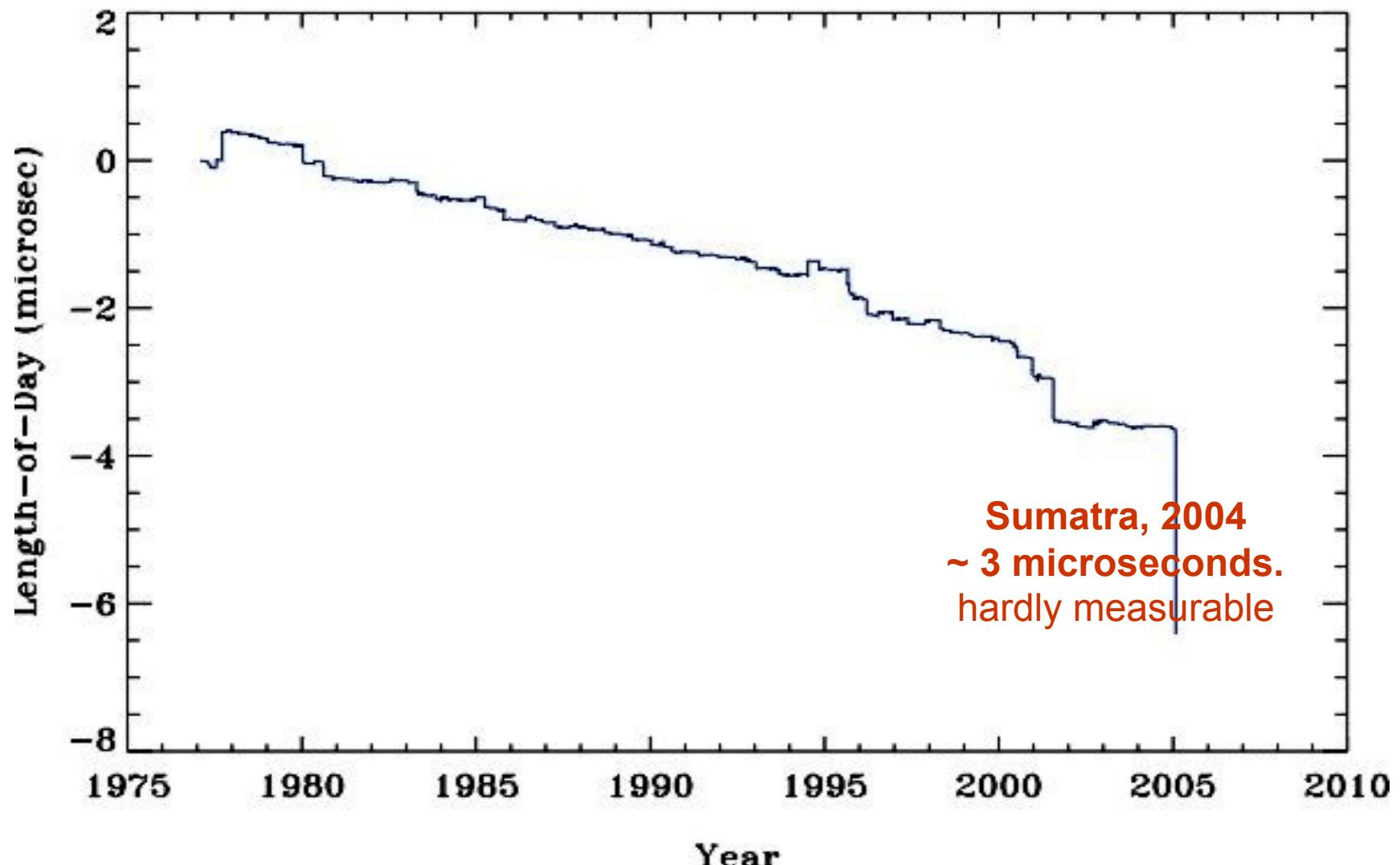
## Co-Seismic Effects on Earth's Rotation

- Milne (1906); Cecchini (1928)
- Munk & MacDonald (1960)
- Alaska earthquake (1964), Press (1965)
- Mansinha & Smylie (1967; +)
- Ben-Menahem & Israel (1970; +)
- Rice & Chinnery (1972)
- Dahlen (1973)
- Dziewonski & O'Connell (1975)
- Smith (1977)
- Souriau & Cazenave (1985)
- Gross (1986)
- Chao & Gross (1987; +)
  - using normal mode summation (Gilbert, 1970)
  - In terms of seismic moment tensor (Harvard CMT catalog)
  - need (SNREI) Earth model (PREM) and normal mode eigen-functions (Masters)
  - similar formulas for changes in gravity field, energy, etc.

# Seismic Moment Tensor

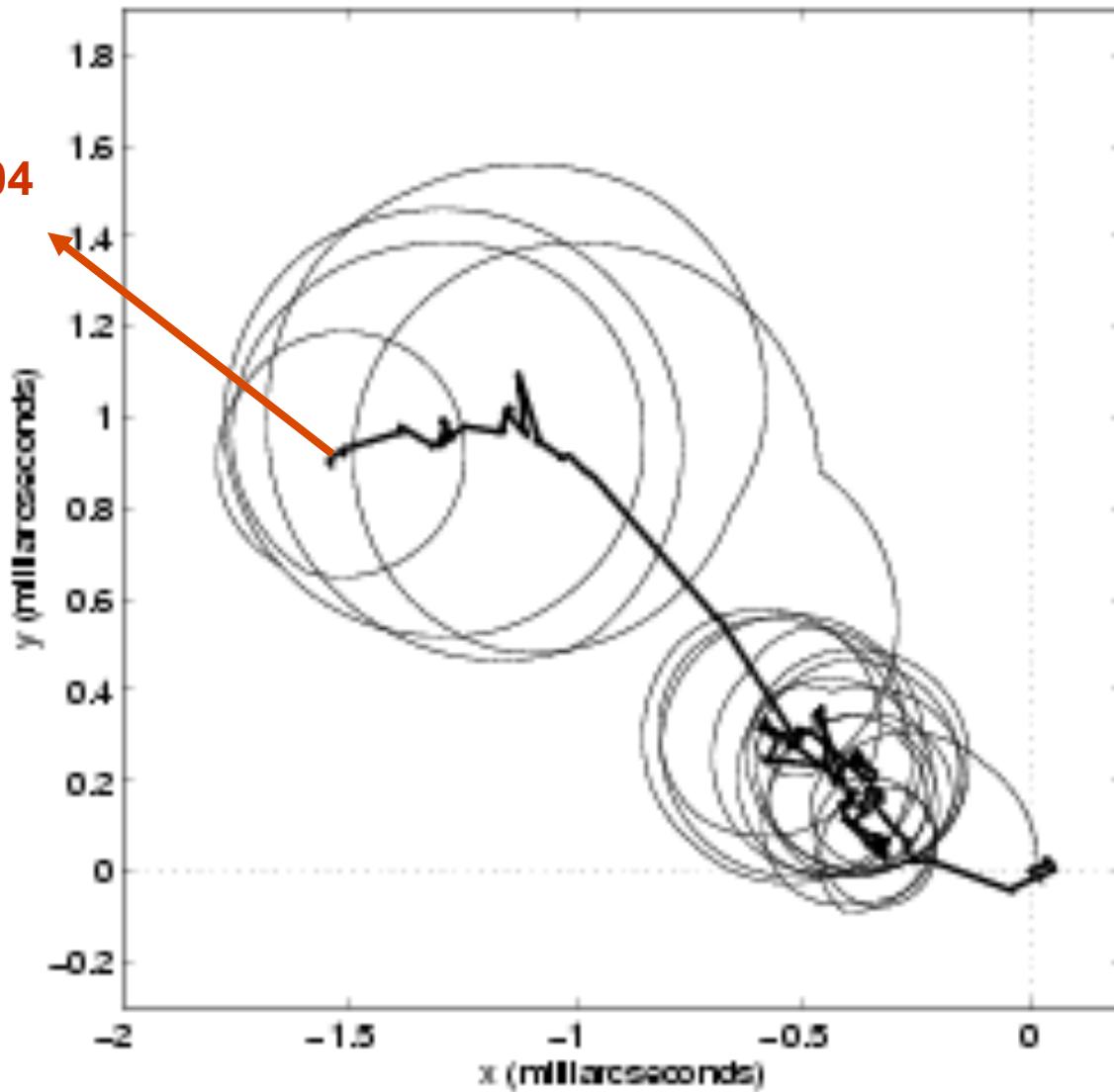
- Contains all information about source mechanism (magnitude, fault direction, slip angles, etc.)
- 2nd-order tensor (conservation of linear momentum)
- Symmetric tensor, only 6 independent parameters (conservation of angular momentum)
- Magnitude (seismic moment)  $[M:M]^{1/2}$  is a good measure of earthquake size
  - => moment scale (vs. Richter scale)

## Cumulative change in Length-of-Day by earthquakes since 1976



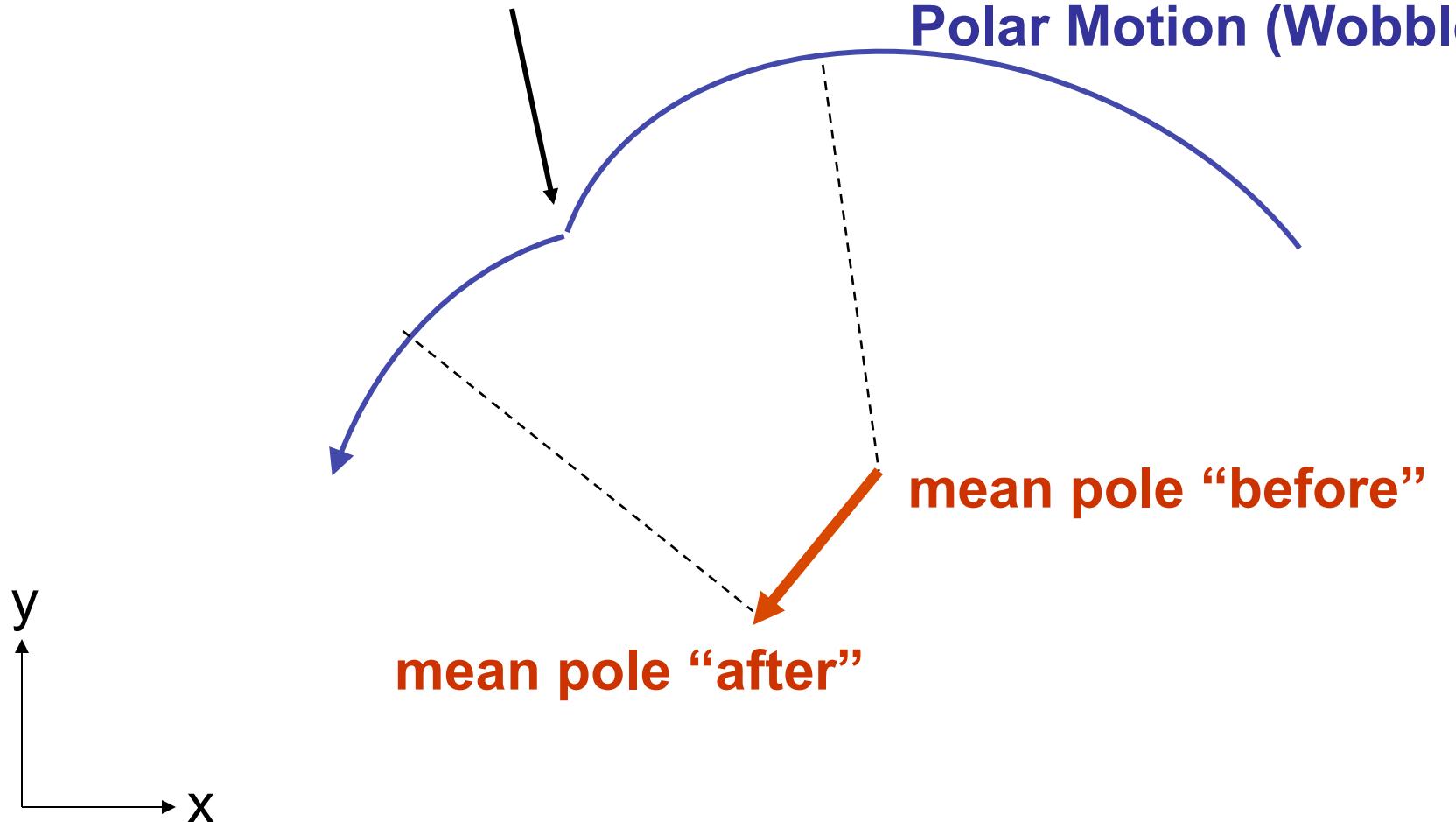
## Cumulative “Mean” Pole Position shift by earthquakes, 1976-1999

**Sumatra, 2004**  
~ 2.5 cm.  
Measurable,  
but “buried”



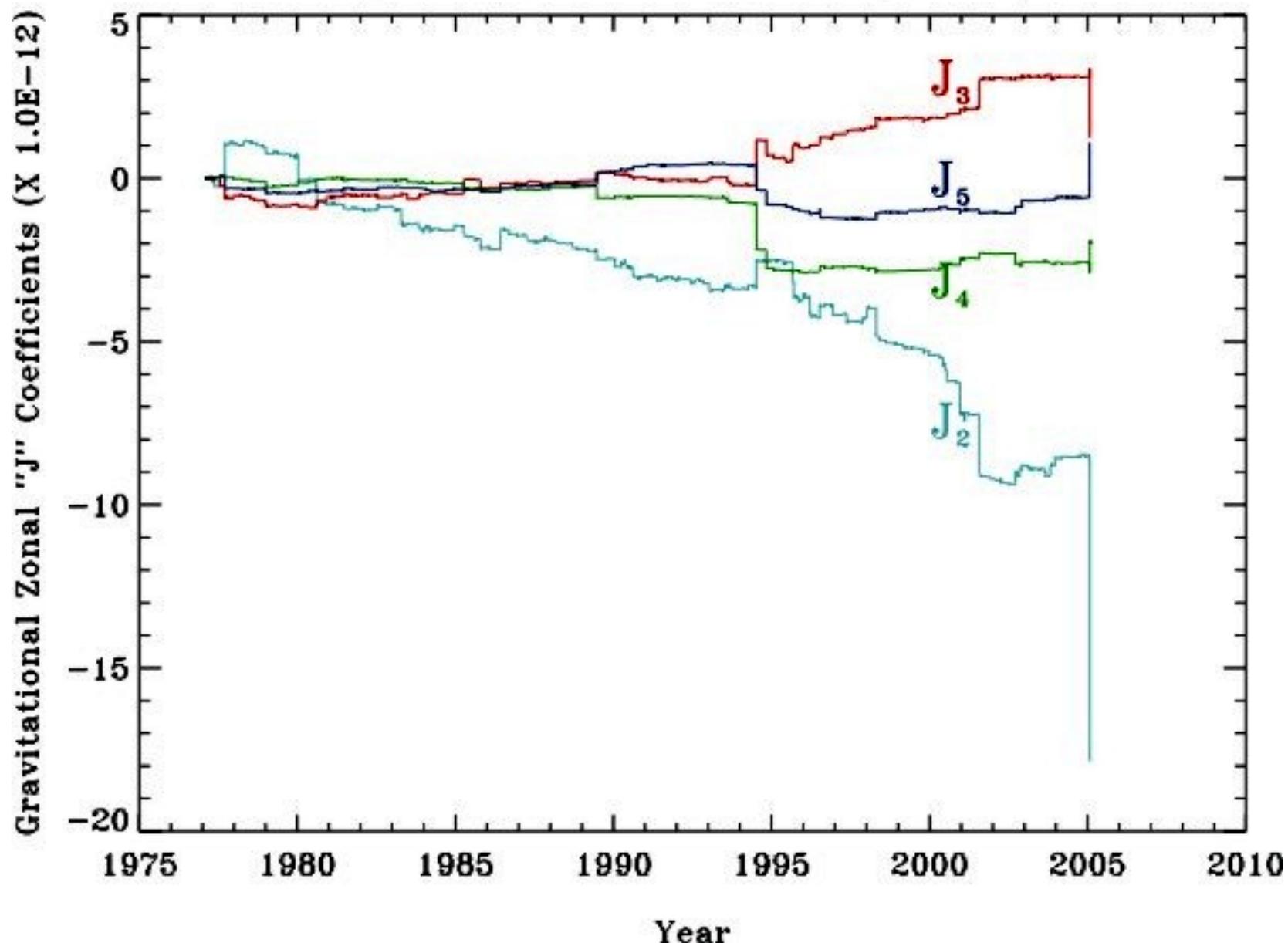
**Earthquake**  
**(step function excitation)**

**Polar Motion (Wobble)**

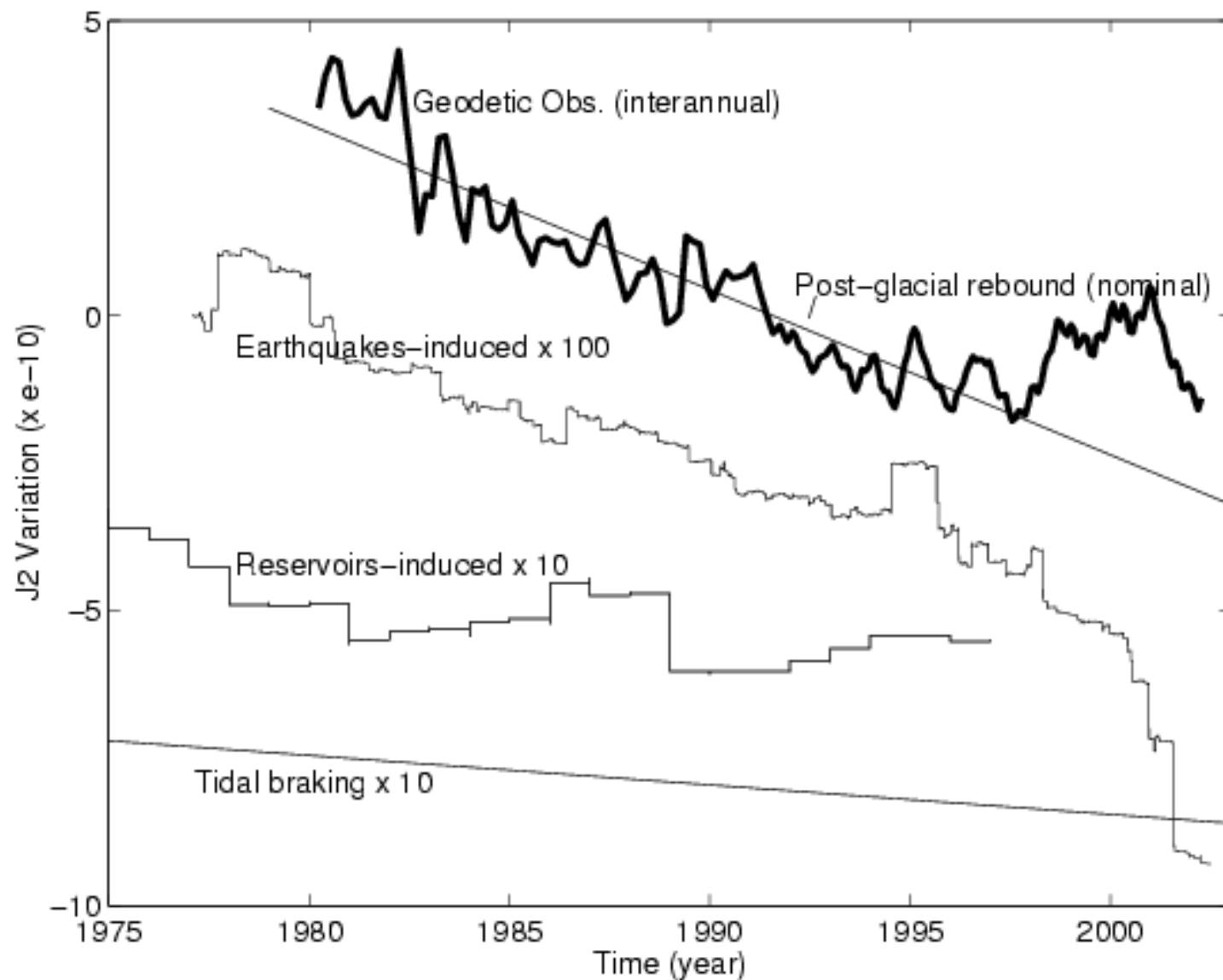


發生時間	地點	地震規模	日長改變量	自轉軸偏移量	自轉軸偏移方向
1957/3/9	阿拉斯加	9.1	(資料不足)	(資料不足)	(資料不足)
1960/5/22	智利	9.5	-8.4 $\mu\text{s}$	68 cm	115 °E
1964/3/27	阿拉斯加	9.2	+ 6.8 $\mu\text{s}$	23 cm	198 °E
2004/12/26	蘇門答臘	9.3	-6.8 $\mu\text{s}$	7 cm	127 °E
2010/2/27	智利	8.8	-1.3 $\mu\text{s}$	8 cm	112 °E
2011/3/11	日本	9.1	-1.6 $\mu\text{s}$	15 cm	139 °E

Cumulative change due to 22369 major earthquakes  
(based on Chao & Gross, 1987)



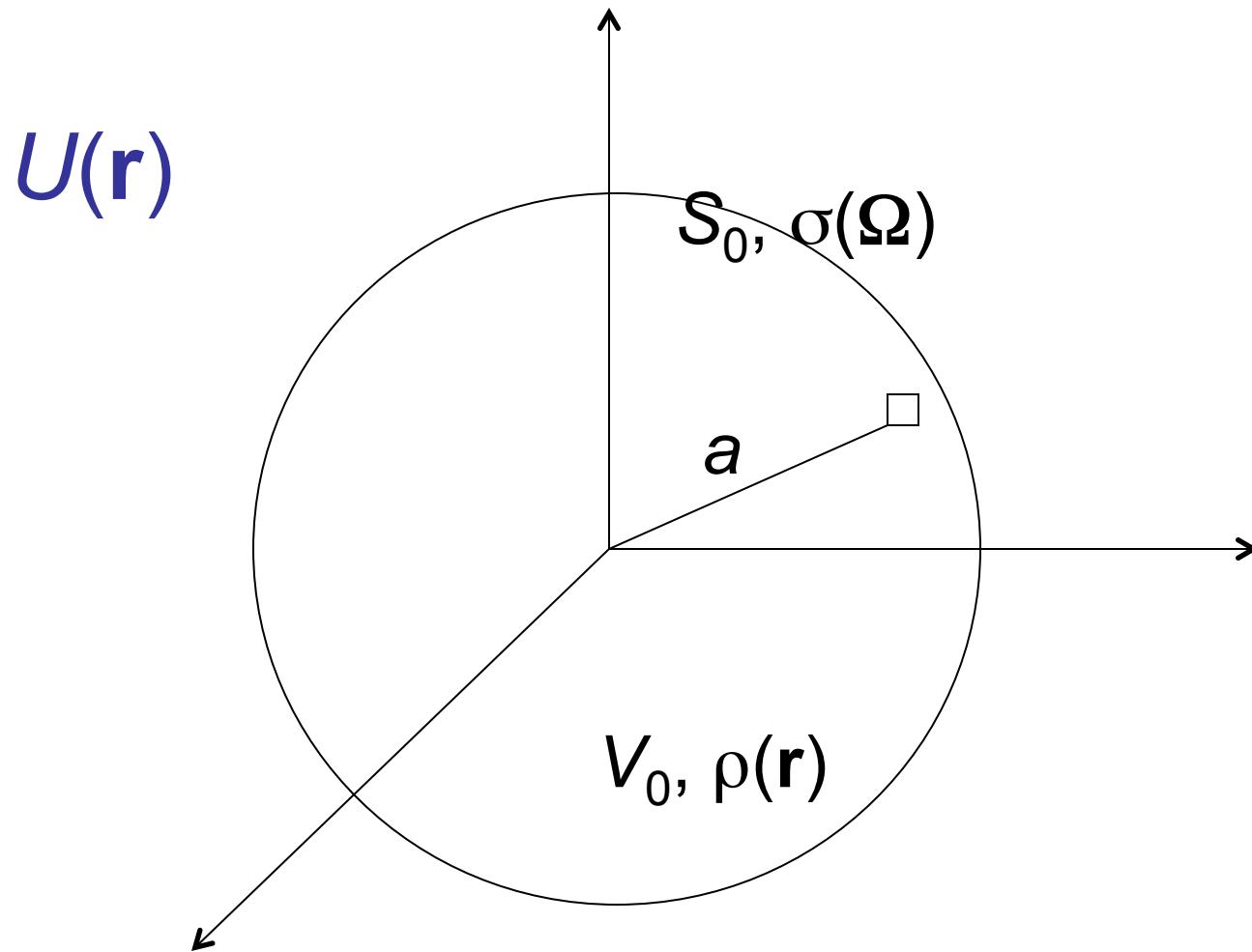
# Long-term changes in Earth's oblateness $J_2$



## Limitation of (time-variable) gravity signal

***“You don’t know where it comes from!”***

- Low spatial resolution
- Non-uniqueness in inversion
- Sum of all sources



## Gravitational Potential Field

- Newton's gravitational law

$$U(\mathbf{r}) = G \iiint_{V_0} \frac{\rho(\mathbf{r}_0)}{|\mathbf{r} - \mathbf{r}_0|} dV_0$$

- Addition theorem

$$\frac{1}{|\mathbf{r} - \mathbf{r}_0|} = \sum_{n=0}^{\infty} \sum_{m=-n}^n (2n+1)^{-1} (r_0^n / r^{n+1}) Y_{nm}^*(\Omega) Y_{nm}(\Omega_0)$$



Multipole expansion of gravity field

$$U(\mathbf{r}) = G \sum_{n=0}^{\infty} \sum_{m=-n}^n \frac{1}{(2n+1)r^{n+1}} \left[ \iiint_{V_0} \rho(\mathbf{r}_0) r_0^n Y_{nm}(\Omega_0) dV_0 \right] Y_{nm}^*(\Omega)$$

## Gravitational Potential Field (Geoid)

- Multipole expansion of Newton's formula:

$$U(\mathbf{r}) = G \sum_{n=0}^{\infty} \sum_{m=-n}^n \frac{1}{(2n+1)r^{n+1}} \left[ \iiint_{V_0} \rho(\mathbf{r}_0) r_0^n Y_{nm}(\Omega_0) dV_0 \right] Y_{nm}^*(\Omega)$$

- Conventional expression (satisfying Laplace Eq. in terms of Stokes Coeff.):

$$U(r, \theta, \lambda) = \frac{GM}{a} \sum_{n=0}^{\infty} \sum_{m=0}^n \left( \frac{a}{r} \right)^{n+1} P_{nm}(\cos \theta) (C_{nm} \cos m\lambda + S_{nm} \sin m\lambda)$$



$$C_{nm} + iS_{nm} = \frac{1}{(2n+1)Md^n} \iiint_{V_0} \rho(\mathbf{r}) r^n Y_{nm}(\Omega) dV$$

# 3-D Gravitational Inversion

- Multipole expansion

$$U(\mathbf{r}) = G \sum_{n=0}^{\infty} \sum_{m=-n}^n \frac{1}{(2n+1)r^{n+1}} \left[ \iiint_{V_0} \rho(\mathbf{r}_0) r_0^n Y_{nm}(\Omega_0) dV_0 \right] Y_{nm}^*(\Omega)$$

$2n + 1$  (known) multipoles for each degree  $n$

- Moment expansion

$$U(\mathbf{r}) = G \sum_{n=0}^{\infty} \sum_{\alpha, \beta, \gamma \geq 0}^{\alpha+\beta+\gamma=n} \frac{(-1)^n}{\alpha! \beta! \gamma!} \left[ \iiint_{V_0} x_0^\alpha y_0^\beta z_0^\gamma \rho(\mathbf{r}_0) dV_0 \right] \frac{\partial^n}{\partial x^\alpha \partial y^\beta \partial z^\gamma} \left( \frac{1}{|\mathbf{r}|} \right)$$

$(n+1)(n+2)/2$  (unknown) moments for each  $n$



“Degree of deficiency” of knowledge is  
 $n(n-1)/2$  for each degree  $n$ .

Degree $n$	# multipoles $(2n + 1)$	# moments $(n+1)(n+2)/2$	Degree of deficiency $n(n-1)/2$
0	1 (monopole)	1 (total mass)	0
1	3 (dipole)	3 (center of mass)	0
2	5 (quadrupole)	6 (inertia tensor)	1
3	7 (octupole)	10 (3rd moment)	3
4	9	15	6
5	11	21	10
6	13	28	15
100	201	5151	4950

The degree of deficiency as a function of spherical harmonic degree  $n$  in the 3-D gravitational inversion.

## Additional physical/mathematical constraints leading to unique solutions:

- minimum shear energy
- maximum entropy of  $\rho$
- minimum norm-2 variance for the lateral distribution
- .....

## 2-D gravitational Inversion on a spherical shell $S_0$

$$C_{nm} + iS_{nm} = \frac{a^2}{(2n+1)M} \iint_{S_0} \sigma(\Omega) Y_{nm}(\Omega) d\Omega$$
$$\sigma(\Omega) = \sum_{n,m} \sigma_{nm} Y_{nm}^*(\Omega)$$



$$\sigma_{nm} = \frac{(2n+1)M}{4\pi a^2} (C_{nm} + iS_{nm})$$

It is possible to mimic ANY external field by means of some proper surface density on  $S_0$ .

## 2-D gravitational Inversion on a spherical shell $S_0$ (cont'd)

For CHANGES due to mass redistribution on  $S_0$  (taking into account of loading effect), in Eulerian description:

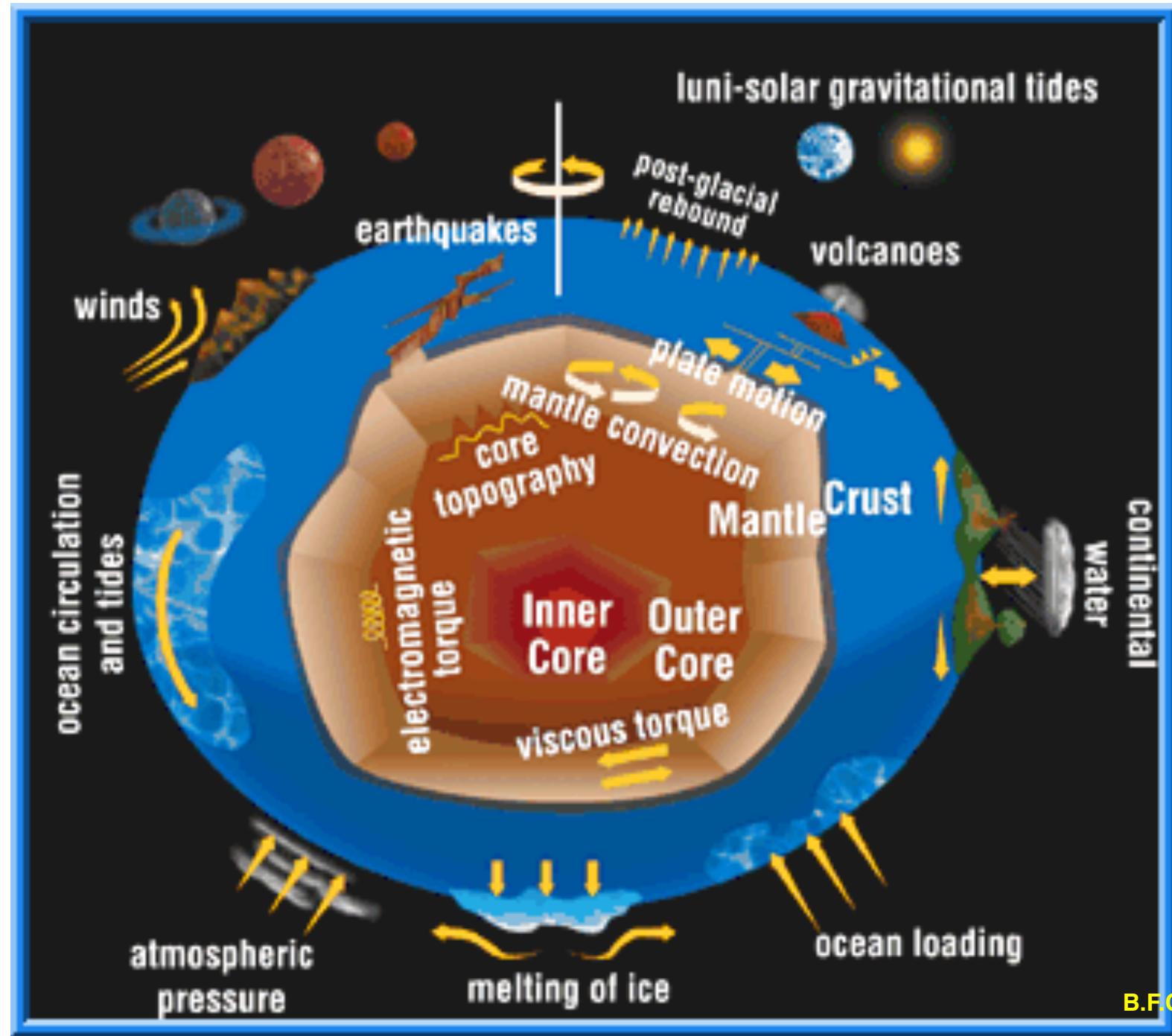
$$\Delta C_{nm}(t) + i \Delta S_{nm}(t) = \frac{a^2(1+k'_n)}{(2n+1)M} \int \int \Delta \sigma(\Omega; t) Y_{nm}(\Omega) d\Omega$$



$$\Delta \sigma_{nm}(t) = \frac{(2n+1)M}{4\pi a^2(1+k'_n)} [\Delta C_{nm}(t) + i \Delta S_{nm}(t)]$$

Unique!

Loading effect “undone”



B.F.Chao

# (near) Surface Mass Transports

- Earth ellipticity  $\sim \frac{1}{2}$  of  $1/300$   $a \sim 10$  km
- Atmosphere scale height  $\sim 10$  km
- Ocean  $< \sim 5$  km
- Land hydrology  $<$  a few km
- Crustal/topography change  $< \sim 30$  km

# Nice things about spherical harmonics:

- Wavelength, or spatial resolution,  
 $\sim 40,000/2N$  km       $\rightarrow$  Concept of spectrum
- Altitude attenuation  $\sim r^{-n+1}$
- Geoid: Stokes Coeff. ( $C_{nm}, S_{nm}$ )
- Gravity Disturbance:  $(n+1) * (C_{nm}, S_{nm})$
- Gravity Anomaly:  $(n-1) * (C_{nm}, S_{nm})$
- Surface Mass Change:  $(2n+1) * (\Delta C_{nm}, \Delta S_{nm})$

## Conclusions

for the [external gravity => mass density] inversion:

- The 3-D inversion is non-unique (well-known).
- This 3-D non-uniqueness is associated with the radial (depth) dimension.
- Comparing the (spherical harmonic) multipole expansion and the moment expansion => The degree of deficiency in inversion is  $n(n-1)/2$  for each degree  $n$ .
- The 2-D inversion on a spherical shell is unique.
- In terms of spherical harmonics this 2-D uniqueness is convenient and useful in (global) time-variable gravity studies (such as GRACE).